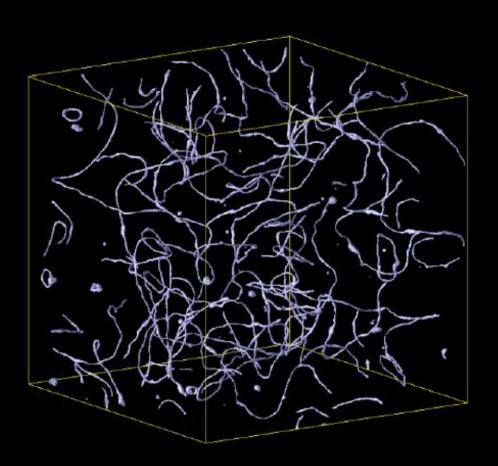
# Y-junction intercommutations of current carrying strings

D.Steer, M.Lilley, DY, T.Hiramatsu, arXiv:1710.07475

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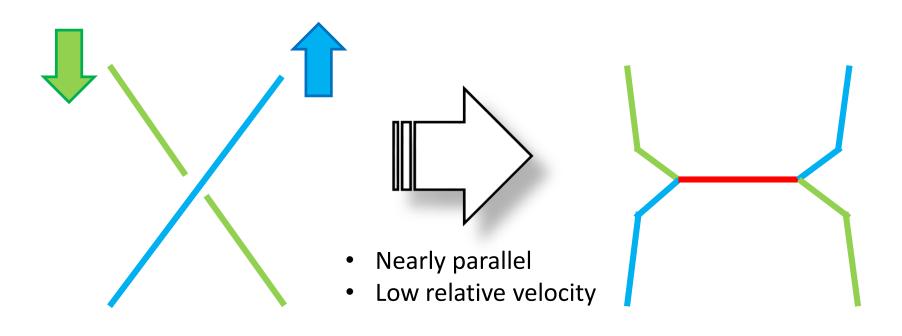


# Cosmic string network

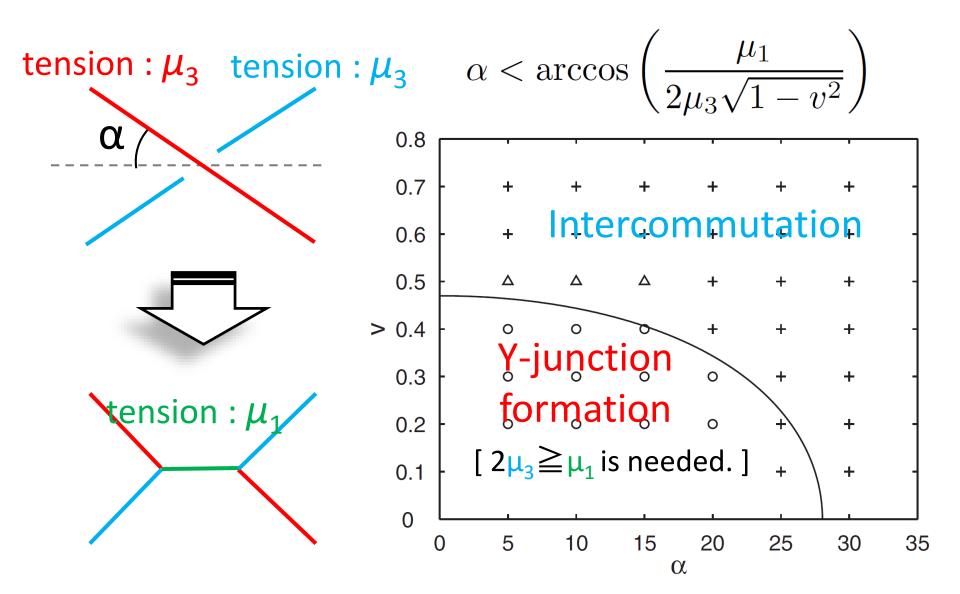
- > Evolves according to scaling.
- Eventually loses its energy through loop formation due to partner exchange: reconnection.
- ➤ Although the reconnection almost always happens, strings with *Y-typed junctions* occurs in many models!

# Y-junction formation?

When two strings collide, bound states and Y-junctions would be produced.



# Y-junction formation: condition



# Strategy [Nambu-Goto case]

(1) Consider effective action

$$S = -\mu \int d\tau d\sigma \sqrt{-\gamma_{ab}}$$

Choose gauge conditions

Junction position at  $\sigma = s_j(\tau)$ 

2 Extend 1 to include multiple strings with junctions

$$S = -\sum_{j} \mu_{j} \int d\tau d\sigma \,\Theta(s_{j}(\tau) - \sigma) \sqrt{-\gamma_{j}}$$

$$+\sum_{i}\int d\tau \mathfrak{f}_{j}\cdot\left[x_{i}(s_{j}(\tau),\tau)-X(\tau)\right]$$

3 Solve junction conditions

4-dim embedding of a junction

#### Questions we are interested in here are:

- What happens to the currents when two current-carrying strings collide?
- Can junctions actually form?

We would like to look for the condition for the junction formation in the analytic way (if possible).

# Differences between Nambu-Goto and current-carrying strings

Conformal and temporal gauges

Worldsheet gauge choices generally made to study NG DO NOT apply to general elastic strings



We need to develop a fully covariant formalism.

Equations of motion are generally NOT integrable

[exception : chiral strings]



We need to solve *junction conditions* for string position and current *simultaneously*.

# Step 1: effective action

 $\succ$  Lagrangian depends on its internal degree of freedom  $\varphi$  :

$$S = \int d^2 \sigma \sqrt{-\det(\gamma_{ab})} \, \mathcal{L}(w)$$

[Carter(1989a,b)]

- $\checkmark$  L is constant → Nambu-Goto strings
- ✓ L is a function of state parameter  $w \rightarrow$  superconducting strings

$$w \equiv \kappa_0 \gamma^{ab} \varphi_{,a} \varphi_{,b}$$

[Witten(1985)]

[For superconducting strings,  $\varphi$  characterizes a phase of scalar field living on a string.]

#### Energy-momentum tensor

Nambu-Goto strings

Conformal and temporal gauges  $T^{\mu\nu} = U u^{\mu}u^{\nu} - T v^{\mu}v^{\nu}$ 

$$U = T \qquad u^{\mu} \propto x^{\mu}, \tau$$

#### **Energy-momentum tensor**

Current-carrying strings

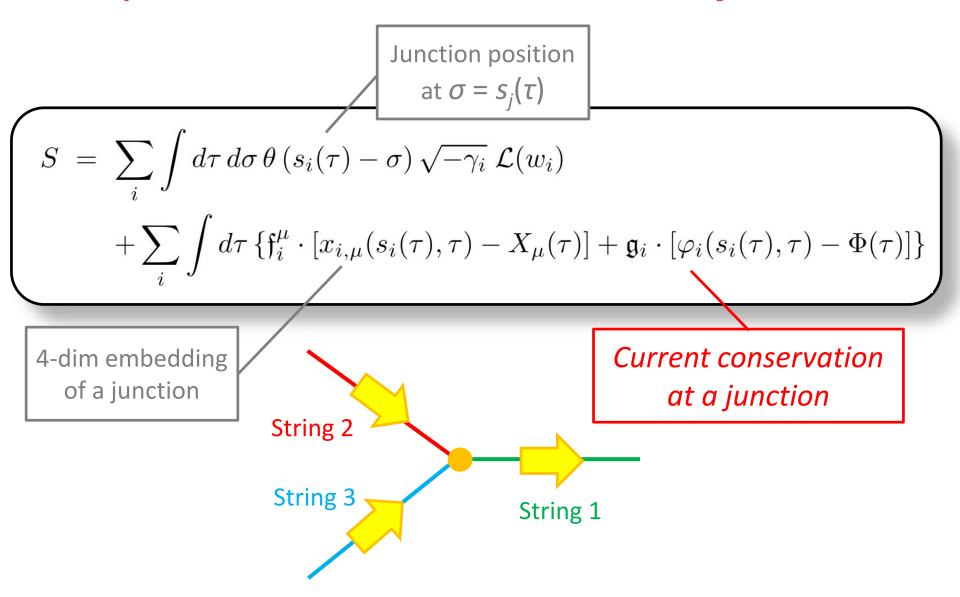
The strings conformal and temporal gauges
$$T^{\mu\nu} = U u^{\mu}u^{\nu} - T v^{\mu}v^{\nu}$$

$$U \neq T \qquad u^{\mu} \propto \varphi^{,a} x^{\mu}_{,a}$$

✓ NOTICE: One can use the residual freedom of Lorentz rotation on the worldsheet!

$$\varphi(\tau,\sigma) \rightarrow \varphi(\sigma)$$
Preferred rest frame

#### Step 2: Effective action with a junction



#### Step 3: Covariant form of junction conditions

NEW

**Current-carrying strings** 

Nambu-Goto strings

Energy-momentum

$$\sum_{j} \Gamma_{j}^{2} \left( U_{j} \dot{s}_{j}^{2} - T_{j} \right) \lambda_{j}^{\mu} = 0$$

$$\sum_{j} \Gamma_j^2 \dot{s}_j \left( U_j - T_j \right) = 0$$

$$\sum_{j} T_{j} \lambda_{j}^{\mu} = 0$$

Outward-directed unit normal vector

$$\sum_{j} \Gamma_{j} \nu_{j} \dot{s}_{j} = 0 \quad [w>0]$$

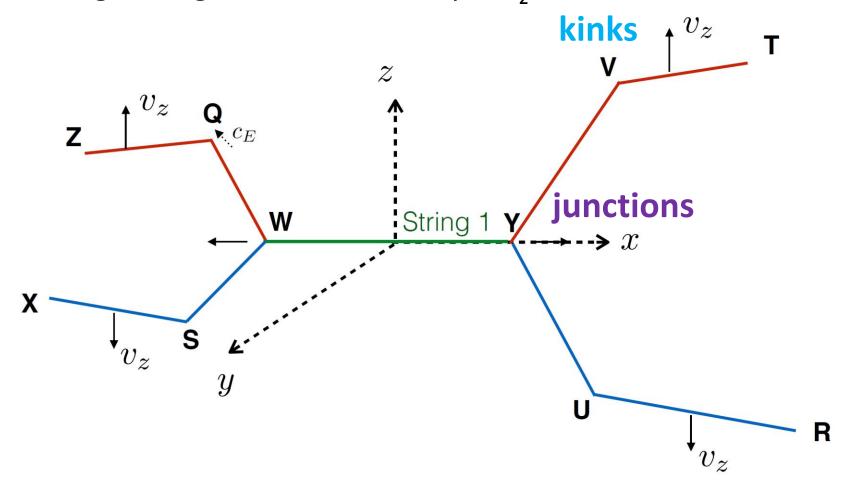
$$\sum_{j} \Gamma_{j} \nu_{j} = 0 \quad [w<0]$$

2 more constraint equations!

Current

# Step 4: Apply to string collision

Consider a string collision between 2 incoming and identical strings at angle  $\pm \alpha$  with velocity  $\pm v_z$ :



# Step 4.1: Nambu-Goto case

> Junction conditions

➤ Unknowns

$$\sum_{j} T_{j} \lambda_{j}^{\mu} = 0 \qquad \Rightarrow \qquad \hat{s}_{3} = \dot{s}_{2} = -\frac{\mu_{3}}{2\mu_{1}} \dot{s}_{1}$$

$$\dot{s}_{1} = \frac{2\mu_{3} \sqrt{1 - v_{z}^{2}} \cos \alpha - \mu_{1}}{2\mu_{3} - \mu_{1} \sqrt{1 - v_{z}^{2}} \cos \alpha}$$

All unknowns can be determined by the junction conditions!



The solution makes sense only if s1>0: the connecting string cannot get shorter.

$$\alpha < \arccos\left(\frac{\mu_1}{2\mu_3\sqrt{1-v^2}}\right)$$

# Step 4.2: Current-carrying case

> Junction conditions

➤ Unknowns

$$\sum_{j} \Gamma_{j}^{2} \left( U_{j} \dot{s}_{j}^{2} - T_{j} \right) \lambda_{j}^{\mu} = 0$$

$$\sum_{j} \Gamma_{j}^{2} \dot{s}_{j} \left( U_{j} - T_{j} \right) = 0$$

$$\sum_{j} \Gamma_{j} \nu_{j} \dot{s}_{j} = 0$$

$$\sum_{j} \Gamma_{j} \nu_{j} \dot{s}_{j} = 0$$

$$\sum_{j} \text{more constraint eqs!}$$

$$\dot{s}_{1} = \frac{r_{here}}{1 - c_{E} v_{+}}$$

$$w_{1} = \cdots \quad w_{n} \quad v_{ariables!}$$

In the presence of the current,  $S_3 = S_2$  is determined **NOT** by the junction conditions but by the configuration.

# Step 4.2: Current-carrying case

Junction conditions

Unknowns

$$\sum_{j} \Gamma_{j}^{2} \left( U_{j} \dot{s}_{j}^{2} - T_{j} \right) \lambda_{j}^{\mu} = 0$$

$$\sum_{j} \Gamma_{j}^{2} \dot{s}_{j} \left( U_{j} - T_{j} \right) = 0$$

$$\sum_{j} \Gamma_{j} \nu_{j} \dot{s}_{j} = 0$$

$$\sum_{j} \Gamma_{j} \nu_{j} \dot{s}_{j} = 0$$

$$\sum_{j} \Gamma_{j} \nu_{j} \dot{s}_{j} = 0$$

$$\sum_{j} \sum_{l} v_{j} \dot{s}_{j} = 0$$

2 more constraint eqs!

(# of eqs to solve) > (# of unknowns)



The system is **OVERDETERMINED**!

## Summary

- ➤ We have extended the analysis of the subsequent formation of Y-junction to the elastic models characterizing current-carrying strings.
- > There are big differences between NG and elastic strings:
  - √ Gauge issues
  - ✓ Presence of internal DoF

In the case of the string collision, the joining string can NOT be described by the elastic model.

## Future prospects

- The treatment of such a collision may generally require the use of a *non-conservative model*.
- ➤ It is of great interest to compare the results from the numerical investigation of a collision of current-carrying strings.

### Thank you!